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三边固支一边自由混凝土矩形薄板的热弯曲

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摘 要: 为了得到温度作用下三边固支一边自由混凝土矩形薄板的挠度和弯矩解析解。该文基于薄板的小挠度理论和叠加原理, 考虑横向变温情况, 将温度作用下的三边固支一边自由矩形薄板看作是面内温差作用下的三边简支一边自由的矩形薄板和相邻三边作用弯矩的三边简支一边自由矩形薄板的叠加。首先, 通过在自由边界上试设具有待定参数的挠度函数, 采用李维解法, 推导出三边简支一边自由矩形薄板在自由边界挠度函数作用下的解析解; 其次, 推导出温度作用下三边简支一边自由矩形薄板的解析解; 再之, 利用相邻三边弯矩作用下四边简支矩形薄板的解答, 推导出相邻三边弯矩作用下三边简支一边自由矩形薄板的解答; 最后, 采用叠加原理得出横向变温作用下三边固支一边自由矩形薄板的挠度和弯矩解析解, 并利用 MATLAB 编制程序得到横向变温作用下三边固支一边自由矩形薄板的计算系数用表, 从而为以后工程结构中三边固支一边自由混凝土矩形薄板在热环境下的设计计算提供了理论依据。

关键词: 混凝土; 三边固支一边自由; 矩形薄板; 温度; 挠度; 弯矩

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THE THERMAL BEND OF CONCRETE RECTANGULAR THIN PLATE WITH THREE CLAMPED SIDES AND ONE FREE SIDE

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Abstract: In order to obtain the deflection and bending analytical solution about concrete rectangular thin plate (RTP) of three clamped sides and one free side, in this paper, based on the small deflection plate theory and superposition principle, considering temperature variation which is perpendicular to surface, the RTP of three clamped sides and one free side under temperature disparity is regarded as under the superposition of temperature disparity and the bending moment with three adjacent sides. Firstly, by supposing deflection function which has undetermined parameter at free side, and adopting Levy method, the closed solution about the RTP under free boundary constraints is obtained; Secondly, the closed solution under temperature disparity is obtained; Thirdly, using the solution of RTP with four simply supported sides under the bending moment on three adjacent sides, the solution of RTP with three simply supported sides and one free side under the bending moment on three adjacent sides is obtained; Finally, the deflection and bending analytical solution about RTP of three clamped sides and one free side under temperature disparity is acquired adopting the superposition principle, and calculation coefficient table about concrete RTP of three clamped sides and one free side under temperature disparity is obtained using MATLAB. And so a theory calculation base about concrete RTP with three clamped sides and one side free under temperature disparity in engineering structure is provided for engineering design.

Key words: concrete; three sides clamped and one side free; rectangular thin plate; temperature; deflection; bending

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矩形薄板广泛应用于土木工程、水利、石油化工、铁路、公路等诸多领域，诸如有顶盖的矩形贮液结构、钢筋混凝土屋面板、剪力墙结构、工业地坪及建筑物的基础、公路路面、机场跑道、停机坪等，这些矩形薄板根据边界条件广义上可分四边支承矩形薄板和具有自由边界的矩形薄板两种情况。在温度作用下，这类结构表面温度迅速上升或下降，由于混凝土材料本身的热惰性，致使结构内部大部分区域仍处于原来的温度状态，从而在薄板的厚度方向形成较大的温差。这种温差作用下的变形受到结构中内部多余约束的制约，从而产生不可忽略的温度应力，从而导致结构构件开裂。课题组历时多年推导了四边简支、四边固支、三边固支一边简支、一边固定三边简支、两邻边固定两邻边简支和两对边固定两对边简支等 6 种四边支承矩形薄板在横向变温作用的挠度和内力解析解^[1-9]，考虑到工程应用方便，利用 MATLAB 软件编制程序计算，得到关于这类混凝土矩形薄板的挠度和内力计算系数用表^[10]。同时，对温度作用下混凝土矩形薄板的热屈曲和热振动展开研究^[11-15]。

对温度作用下具有自由边界的矩形薄板，有关文献^[16]虽列出了三边固定一边自由矩形薄板的计算用表(见表 1)，但固定边横向弯矩均大于由 $w=0$ 所得出的解答，对于温度作用下其它情况的矩形薄板，目前国内外文献尚未见相关报导。因此，本文首先针对三边固支一边自由的矩形薄板，基于薄板

表 1 温差作用下三边固定一边自由矩形薄板的弯矩系数
Table 1 Bending calculation coefficient of three clamped sides and one free side under temperature disparity

$$\mu = \frac{1}{6}, M_x^T = k_x^T \alpha \Delta T E h^2$$

$$M_y^T = k_y^T \alpha \Delta T E h^2$$

$$M_{x1}^{OT} = k_{x1} \alpha \Delta T E h^2$$

$$M_{y1}^{OT} = k_{y1} \alpha \Delta T E h^2$$

$$M_{x2}^{OT} = k_{x2} \alpha \Delta T E h^2$$

$$M_{y2}^{OT} = k_{y2} \alpha \Delta T E h^2$$

$$T = \frac{T_2 + T_1}{2} - \frac{(T_2 - T_1)}{h} z$$

l_x/l_y	k_{x1}	k_{y1}	k_{x2}	k_{y2}	k_x	k_y
0.50	0.1018	0.0983	0.0973	0.0975	0.0948	0.0974
0.75	0.1057	0.0980	0.0973	0.1004	0.0925	0.0913
1.00	0.1085	0.0968	0.0974	0.1050	0.0919	0.0851
1.25	0.1072	0.0957	0.0979	0.1085	0.0931	0.0768
1.50	0.1006	0.0965	0.0983	0.1091	0.0951	0.0696
1.75	0.0997	0.0943	0.0975	0.1013	0.0969	0.0633
2.00	0.0981	0.0933	0.0963	0.0957	0.0985	0.0570

的小挠度理论和叠加原理，推导横向变温作用下三边固支一边自由矩形薄板的解析解，并就混凝土材料利用 MATLAB 软件编制程序得到计算系数用表，为以后工程结构中三边固支一边自由混凝土矩形薄板在热环境下的设计计算提供理论依据。

1 薄板热弹性问题的基本方程

1.1 计算假定

- 1) 变形前垂直于中面的直线，在薄板变形后仍然垂直于变形后的中面，且长度保持不变。
- 2) 应力分量(σ_z 、 τ_{xz} 、 τ_{yz})远小于其余三个应力分量(σ_x 、 σ_y 、 τ_{xy})，因而它们引起的应变可忽略不计。
- 3) 薄板中面内的各点都没有平行于中面的位移， $u|_{z=0}=0$ ， $v|_{z=0}=0$ 。
- 4) 忽略混凝土浇注过程中所产生的热应力。

1.2 热弹性问题基本方程的推导

薄板内任一点的应力分量为：

$$\begin{cases} \sigma_x = -\frac{Ez}{1-\mu^2} \left(\frac{\partial^2 w}{\partial x^2} + \mu \frac{\partial^2 w}{\partial y^2} \right) - \frac{E\alpha T}{1-\mu} \\ \sigma_y = -\frac{Ez}{1-\mu^2} \left(\frac{\partial^2 w}{\partial y^2} + \mu \frac{\partial^2 w}{\partial x^2} \right) - \frac{E\alpha T}{1-\mu} \\ \tau_{xy} = \tau_{yx} = -\frac{Ez}{1+\mu} \frac{\partial^2 w}{\partial x \partial y} \end{cases} \quad (1)$$

由于薄板厚度与另外两个方向的尺寸相比很小，故假定温度只沿厚度方向变化；另外，由于是薄板，在板结构正常使用之前，混凝土凝结硬化时的发热随着浇注时间的变化而趋向于零，这样原抛物线的非线性温度分布规律 T 变为沿横向线性变化的情形，即：

$$T = \frac{T_2 + T_1}{2} - \frac{(T_2 - T_1)}{h} z \quad (2)$$

式中， T_1 和 T_2 分别为薄板上下表面的温度。

根据现有文献^[17]，热载作用下混凝土的弹性模量 $E(T)$ 与常温下混凝土的弹性模量 E 的关系可由下式确定：

$$\begin{cases} E(T) = E, & T \leq 60 \\ E(T) = 0.88E \sim 0.94E, & 60 < T \leq 100 \\ E(T) = 0.95E \sim 1.08E, & 100 < T \leq 300 \\ E(T) = \left[1 + 18 \left(\frac{T}{1000} \right)^{5.1} \right]^{-1} E, & T > 300 \end{cases}$$

温度作用下线膨胀系数 $\alpha(T)$ 可由下式确定：

$$\alpha(T) = 28 \left(\frac{T}{1000} \right) \times 10^{-6}$$

可以看出,对于常温情况,温度对混凝土的弹性模量和线膨胀系数的影响不大,故近似取:

$$E(T) = E, \quad \alpha(T) = \alpha$$

因此,如图1所示,设 M_x 、 M_y 为横截面单位宽度上的弯矩, M_{xy} 为横截面单位宽度上的扭矩,

F_{Qy} 为横向剪力,并令 $M^T = \frac{E\alpha\Delta Th^2}{12(1-\mu)}$, 则有:

$$\begin{cases} M_x = -D \left(\frac{\partial^2 w}{\partial x^2} + \mu \frac{\partial^2 w}{\partial y^2} \right) - M^T \\ M_y = -D \left(\frac{\partial^2 w}{\partial y^2} + \mu \frac{\partial^2 w}{\partial x^2} \right) - M^T \\ M_{xy} = -D(1-\mu) \frac{\partial^2 w}{\partial x \partial y} \\ F_{Qy} = -D \frac{\partial}{\partial y} \nabla^2 w \end{cases} \quad (3)$$

由于等厚度薄板弹性曲面的平衡微分方程为:

$$\frac{\partial^2 M_x}{\partial x^2} + 2 \frac{\partial^2 M_{xy}}{\partial x \partial y} + \frac{\partial^2 M_y}{\partial y^2} + q = 0$$

故只有温度没有外荷载,故将式(3)代入上式,有:

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right) w = 0 \quad (4)$$

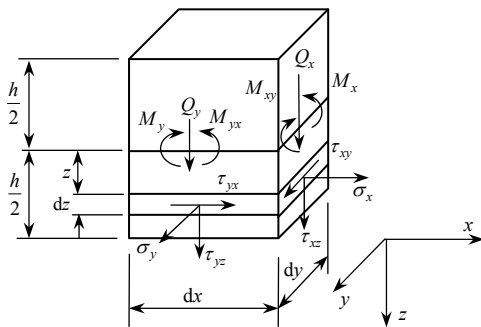


图1 单元内力及应力示意图

Fig.1 Sketch map of element internal force stress

2 热弹性问题的解析解

2.1 边界条件

如图2,对固定边,边界条件为:

$$w \Big|_{x=0} = 0, \quad \frac{\partial w}{\partial x} \Big|_{x=0} = 0 \quad (5)$$

$$w \Big|_{y=0} = 0, \quad \frac{\partial w}{\partial y} \Big|_{y=0} = 0 \quad (6)$$

对自由边,薄板的弯矩 M_y 、扭矩 M_{yx} 以及横向剪力 F_{Qy} 都等于零,并设在 $y=b$ 的边界上的挠度由正弦级数来表示,故有:

$$\begin{cases} M_y \Big|_{y=b} = 0; & M_{yx} \Big|_{y=b} = 0 \\ F_{Qy} \Big|_{y=b} = 0 \end{cases} \quad (7)$$

$$w \Big|_{y=b} = \sum_{m=1,3,\dots} a_m \sin \frac{m\pi x}{a} \quad (8)$$

由式(3)、式(7)的第一式变为:

$$\frac{\partial^2 w}{\partial y^2} + \mu \frac{\partial^2 w}{\partial x^2} \Big|_{y=b} = -\frac{M^T}{D} \quad (9)$$

设自由边界上的分布剪力为 \bar{F}_{Qy} , 则式(7)的第二式和第三式合并为:

$$\bar{F}_{Qy} \Big|_{y=b} = -D \left[\frac{\partial^3 w}{\partial y^3} + (2-\mu) \frac{\partial^3 w}{\partial x^2 \partial y} \right] \Big|_{y=b} = 0 \quad (10)$$

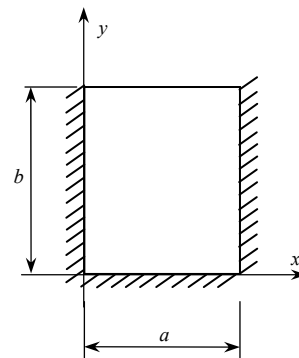


图2 三边固支一边自由矩形薄板

Fig.2 Three clamped sides and one free side

2.2 热弹性问题的解析解

可以看出,要寻找满足完全边界条件式(5)、式(6)、式(8)、式(9)和式(10)的挠度函数非常困难,故采用叠加原理解决这一问题。

由于在固定边界上 $w=0$, 根据式(3)、式(5)和式(6)可知:

$$M_x \Big|_{x=0} = -M^T, \quad M_y \Big|_{y=0} = -M^T$$

可见横向变温作用下的三边固定一边自由矩形薄板可看作温差 ΔT 作用下的三边简支一边自由的矩形薄板与相邻三边作用弯矩 M^T 的三边简支一边自由矩形薄板的叠加。

2.2.1 挠度 $w \Big|_{y=b}$ 作用下三边简支一边自由矩形薄板的解答

薄板的解答

由边界条件式(8), 有:

$$w|_{y=b} = \sum_{m=1,3,\dots} a_m \sin \frac{m\pi x}{a} \quad (11)$$

$$w|_{x=0} = 0, \quad \frac{\partial^2 w}{\partial x^2} \Big|_{x=0} = 0 \quad (12)$$

$$w|_{y=0} = 0, \quad \frac{\partial^2 w}{\partial y^2} \Big|_{y=0} = 0 \quad (13)$$

$$\frac{\partial^2 w}{\partial y^2} + \mu \frac{\partial^2 w}{\partial x^2} \Big|_{y=b} = 0 \quad (14)$$

由式(4), 令:

$$w = -\sum_{m=1}^{\infty} X_m Y_m$$

按照边界条件(10), 令 $X_m = \sin \frac{m\pi x}{a}$, 因此挠度函数 w 可写作为:

$$w = -\sum_{m=1}^{\infty} Y_m \sin \frac{m\pi x}{a} \quad (15)$$

将式(15)代入微分方程式(4)得:

$$Y_m^{(4)} - 2 \frac{m^2 \pi^2}{a^2} Y_m'' + \frac{m^4 \pi^4}{a^4} Y_m = 0$$

这个方程的解可写作如下形式:

$$Y_m = A_m \sinh \frac{m\pi y}{a} + B_m \cosh \frac{m\pi y}{a} + C_m \frac{m\pi y}{a} \sinh \frac{m\pi y}{a} + D_m \frac{m\pi y}{a} \cosh \frac{m\pi y}{a}$$

即有:

$$\begin{cases} M_x = -\frac{D}{2a^2} (1-\mu)^2 \pi^2 \times \\ \sum_{m=1,3,\dots}^{\infty} \left\{ \frac{a_m m^2}{\sinh \beta_m} \sin \frac{m\pi x}{a} \times \left[\left(\beta_m \coth \beta_m + 2 \frac{1+\mu}{1-\mu} \right) \sinh \frac{m\pi y}{a} - \frac{m\pi y}{a} \cosh \frac{m\pi y}{a} \right] \right\} \\ M_y = -\frac{D}{2a^2} (1-\mu)^2 \pi^2 \times \sum_{m=1,3,\dots}^{\infty} \frac{a_m m^2}{\sinh \beta_m} \left(\frac{m\pi y}{a} \cosh \frac{m\pi y}{a} - \beta_m \coth \beta_m \sinh \frac{m\pi y}{a} \right) \sin \frac{m\pi x}{a} \end{cases} \quad (19)$$

2.2.2 ΔT 作用下三边简支一边自由矩形薄板的解析解

1) 边界条件。

如图 3, 边界条件为:

$$w_1|_{x=0} = 0; \quad M_x|_{x=0} = 0 \quad (20)$$

$$w_1|_{y=0} = 0; \quad M_y|_{y=0} = 0 \quad (21)$$

$$\begin{cases} M_y|_{y=0} = 0; \quad M_{xy}|_{y=b} = 0 \\ F_{Qy}|_{y=b} = 0 \end{cases} \quad (22)$$

$$w_1 = \sum_{m=1}^{\infty} \left(A_m \sinh \frac{m\pi y}{a} + B_m \cosh \frac{m\pi y}{a} + C_m \frac{m\pi y}{a} \sinh \frac{m\pi y}{a} + D_m \frac{m\pi y}{a} \cosh \frac{m\pi y}{a} \right) \sin \frac{m\pi x}{a} \quad (16)$$

将其代入式(12), 边界条件自然满足, 将其代入式(13)得:

$$B_m = C_m = 0$$

将其代入式(11)和式(14), 并令 $\frac{m\pi b}{a} = \beta_m$ 得:

$$\begin{cases} A_m \sinh \beta_m + D_m \beta_m \cosh \beta_m = a_m \\ A_m (1-\mu) \sinh \beta_m + \\ D_m [(1-\mu) \beta_m \cosh \beta_m + 2 \sinh \beta_m] = 0 \end{cases} \quad (17)$$

由此可得:

$$\begin{cases} D_m = (\mu-1) \frac{a_m}{2 \sinh \beta_m} \\ A_m = \frac{a_m}{2 \sinh \beta_m} [2 + (1-\mu) \beta_m \coth \beta_m] \end{cases}$$

所以, 三边简支一边自由矩形薄板的挠度表达式为:

$$w = -\sum_{m=1,3,\dots}^{\infty} \left\{ \frac{a_m (1-\mu)}{2 \sinh \beta_m} \times \left[\left(\frac{2}{1-\mu} + \beta_m \coth \beta_m \right) \times \sinh \frac{m\pi y}{a} - \frac{m\pi y}{a} \cosh \frac{m\pi y}{a} \right] \sin \frac{m\pi x}{a} \right\} \quad (18)$$

令 M^T 为零, 将式(18)代入式(3)得:

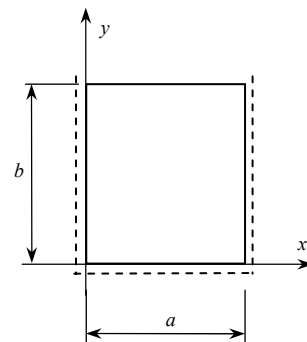


图 3 三边简支一边自由矩形薄板
Fig.3 Three simply supported sides and one free side

对简支边，因挠度 w_1 沿整个边界上都等于零，所以由式(3)，式(20)~式(22)变为：

$$w_1 \Big|_{x=0} = 0 ; \quad \frac{\partial^2 w_1}{\partial x^2} \Big|_{x=0} = -\frac{M^T}{D} \quad (23)$$

$$w_1 \Big|_{y=0} = 0 ; \quad \frac{\partial^2 w_1}{\partial y^2} \Big|_{y=0} = -\frac{M^T}{D} \quad (24)$$

$$\begin{cases} \frac{\partial^2 w_1}{\partial y^2} + \mu \frac{\partial^2 w_1}{\partial x^2} \Big|_{y=b} = 0 \\ \frac{\partial^3 w_1}{\partial y^3} + (2-\mu) \frac{\partial^3 w_1}{\partial x^2 \partial y} \Big|_{y=b} = 0 \end{cases} \quad (25)$$

2) ΔT 作用下的解析解。

由边界条件式(8)，温差 ΔT 作用下的三边简支一边自由的矩形薄板可看作四边简支矩形薄板和边界上作用挠度 $w_1 \Big|_{y=b}$ 的三边简支一边自由矩形薄板的叠加。

温度作用下四边简支矩形薄板的解答为^[2 3]：

$$w_1 = -\frac{4a^2 M^T}{D\pi^3} \times \sum_{m=1,3,\dots} \left\{ \frac{1}{m^3 \cosh \alpha_m} \times \cosh \left(\frac{2\alpha_m y}{b} - \alpha_m \right) \sin \frac{m\pi x}{a} \right\} - \frac{M^T}{2D} (x-a)x \quad (26)$$

$$\begin{cases} M_{x1} = \frac{4M^T}{\pi} (\mu-1) \times \sum_{m=1,3,\dots} \left\{ \frac{1}{m \cosh \alpha_m} \times \cosh \left(\frac{m\pi y}{a} - \alpha_m \right) \sin \frac{m\pi x}{a} \right\} \\ M_{y1} = (\mu-1)M^T + \frac{4M^T}{\pi} (1-\mu) \times \sum_{m=1,3,\dots} \left\{ \frac{1}{m \cosh \alpha_m} \times \cosh \left(\frac{m\pi y}{a} - \alpha_m \right) \sin \frac{m\pi x}{a} \right\} \end{cases} \quad (27)$$

$$M_{x2} = \frac{2(3-2\mu)M^T}{\pi}$$

$$\sum_{m=1,3,\dots} \left\{ \frac{\sinh \beta_m \sinh \left(\frac{2\alpha_m y}{b} - \alpha_m \right)}{m \cosh \alpha_m} \left[\left(\beta_m \coth \beta_m + 2 \frac{1+\mu}{1-\mu} \right) \times \sinh \frac{m\pi y}{a} - \frac{m\pi y}{a} \cosh \frac{m\pi y}{a} \right] \sin \frac{m\pi x}{a} \right\}$$

令 $a_m = \bar{a}_m$ ，将式(18)代入式(10)，得：

$$\bar{F}_{Qy}^1 \Big|_{y=b} = -\frac{D\pi^3(1-\mu)^2}{2a^3} \times \sum_{m=1,3,\dots} \frac{\bar{a}_m m^3}{\sinh^2 \beta_m} \left[\frac{3+\mu}{1-\mu} \frac{\sinh(2\beta_m)}{2} + \beta_m \right] \sin \frac{m\pi x}{a}$$

将式(26)代入式(10)，得：

$$\bar{F}_{Qy}^2 \Big|_{y=b} = -\frac{4(3-2\mu)M^T}{\pi b} \times \sum_{m=1,3,\dots} \frac{\alpha_m}{m \cosh \alpha_m} \sinh \left(\frac{2\alpha_m y}{b} - \alpha_m \right) \sin \frac{m\pi x}{a}$$

在自由边界上，由式(10)，得：

$$\bar{F}_{Qy}^1 \Big|_{y=b} + \bar{F}_{Qy}^2 \Big|_{y=b} = 0$$

即有：

$$\bar{a}_m = -\frac{8a^3(3-2\mu)M^T \sinh^2 \beta_m}{bD\pi^4(1-\mu)^2} \times \frac{\alpha_m \sinh \left(\frac{2\alpha_m y}{b} - \alpha_m \right)}{m^4 \cosh \alpha_m} \frac{3+\mu}{1-\mu} \frac{\sinh(2\beta_m)}{2} + \beta_m$$

将其代入式(18)和式(19)得：

$$w_2 = \frac{2a^2(3-2\mu)M^T}{D\pi^3(1-\mu)} \times \sum_{m=1,3,\dots} \left\{ \frac{\sinh \beta_m \sinh \left(\frac{2\alpha_m y}{b} - \alpha_m \right)}{m^3 \cosh \alpha_m} \frac{3+\mu}{1-\mu} \frac{\sinh(2\beta_m)}{2} + \beta_m \right\} \left[\left(\frac{2}{1-\mu} + \beta_m \coth \beta_m \right) \cdot \sinh \frac{m\pi y}{a} - \frac{m\pi y}{a} \cosh \frac{m\pi y}{a} \right] \sin \frac{m\pi x}{a} \quad (28)$$

$$M_{y2} = \frac{2(3-2\mu)M^T}{\pi} \sum_{m=1,3,\dots}^{\infty} \left\{ \frac{\sinh \beta_m \sinh \left(\frac{2\alpha_m y}{b} - \alpha_m \right)}{m \cosh \alpha_m} \left(\frac{m\pi y}{a} \cosh \frac{m\pi y}{a} - \beta_m \coth \beta_m \sinh \frac{m\pi y}{a} \right) \sin \frac{m\pi x}{a} \right\}. \quad (29)$$

2.2.3 相邻三边作用弯矩 M_-^T 的三边简支一边自由矩形薄板的解答

由边界条件式(8), 相邻三边作用弯矩 M_-^T 的三边简支一边自由的矩形薄板可看作相邻三边作用弯矩 M_-^T 的四边简支矩形薄板和在上边界上作用挠度 $w|_{y=b}$ 的三边简支一边自由矩形薄板的叠加。

相邻三边作用弯矩 M_-^T 的四边简支矩形薄板的挠度和弯矩表达式如下^[5]:

$$w_3(x, y) = \frac{8M_-^T}{\pi^4 D} \sum_{i=1,3,\dots}^{\infty} \sum_{j=1,3,\dots}^{\infty} \left\{ \left[\frac{1}{ij} \left(\frac{i^2}{a^2} + \frac{j^2}{b^2} \right)^{-2} \times \left(\frac{2i^2}{a^2} + \frac{j^2}{b^2} \right) \right] \sin \frac{i\pi x}{a} \sin \frac{j\pi y}{b} \right\} \quad (30)$$

$$M_{x3} = \frac{8M_-^T}{\pi^2} \times \sum_{i=1,3,\dots}^{\infty} \sum_{j=1,3,\dots}^{\infty} \left\{ \frac{1}{ij} \left(\frac{i^2}{a^2} + \mu \frac{j^2}{b^2} \right) \times \left(\frac{i^2}{a^2} + \frac{j^2}{b^2} \right)^{-2} \times \left(\frac{2i^2}{a^2} + \frac{j^2}{b^2} \right) \sin \frac{i\pi x}{a} \sin \frac{j\pi y}{b} \right\} \quad (31a)$$

$$\bar{a}_m = -\frac{16M_-^T a^3 \sinh^2 \beta_m}{\pi^4 D b (1-\mu)^2} \times \frac{\sum_{k=1,3,\dots}^{\infty} \left\{ \frac{1}{m^4} \left(\frac{m^2}{a^2} + \frac{k^2}{b^2} \right)^{-2} \left(\frac{2m^2}{a^2} + \frac{k^2}{b^2} \right) \times \left[\frac{k^2}{b^2} + (2-\mu) \frac{m^2}{a^2} \right] \right\}}{\frac{3+\mu}{1-\mu} \frac{\sinh(2\beta_m)}{2} + \beta_m}$$

将其代入式(18)、式(19), 得:

$$w_4 = -\frac{8M_-^T a^3}{\pi^4 D b (1-\mu)} \sum_{m=1,3,\dots}^{\infty} \left\{ \frac{\sinh \beta_m \sum_{k=1,3,\dots}^{\infty} \left\{ \frac{1}{m^4} \left(\frac{m^2}{a^2} + \frac{k^2}{b^2} \right)^{-2} \left(\frac{2m^2}{a^2} + \frac{k^2}{b^2} \right) \left[\frac{k^2}{b^2} + (2-\mu) \frac{m^2}{a^2} \right] \right\}}{\frac{3+\mu}{1-\mu} \frac{\sinh(2\beta_m)}{2} + \beta_m} \times \left[\left(\frac{2}{1-\mu} + \beta_m \coth \beta_m \right) \sinh \frac{m\pi y}{a} - \frac{m\pi y}{a} \cosh \frac{m\pi y}{a} \right] \sin \frac{m\pi x}{a} \right\} \quad (32)$$

$$M_{y3} = \frac{8M_-^T}{\pi^2} \times \sum_{i=1,3,\dots}^{\infty} \sum_{j=1,3,\dots}^{\infty} \left\{ \frac{1}{ij} \left(\mu \frac{i^2}{a^2} + \frac{j^2}{b^2} \right) \times \left(\frac{i^2}{a^2} + \frac{j^2}{b^2} \right)^{-2} \times \left(\frac{2i^2}{a^2} + \frac{j^2}{b^2} \right) \sin \frac{i\pi x}{a} \sin \frac{j\pi y}{b} \right\} \quad (31b)$$

将式(30)代入式(10), 得:

$$\bar{F}_{Qy}^3 \Big|_{y=b} = -\frac{8M_-^T}{\pi b} \times \sum_{i=1,3,\dots}^{\infty} \sum_{j=1,3,\dots}^{\infty} \left\{ \frac{1}{i} \left(\frac{i^2}{a^2} + \frac{j^2}{b^2} \right)^{-2} \left(\frac{2i^2}{a^2} + \frac{j^2}{b^2} \right) \times \left[\frac{j^2}{b^2} + (2-\mu) \frac{i^2}{a^2} \right] \sin \frac{i\pi x}{a} \right\}$$

令 $a_m = \bar{a}_m$, 在自由边界上由式(10), 得:

$$\bar{F}_{Qy}^1 \Big|_{y=b} + \bar{F}_{Qy}^3 \Big|_{y=b} = 0$$

即有:

$$M_{x4} = -\frac{8M^T a}{\pi^2 b} \sum_{m=1,3,\dots}^{\infty} \left\{ \frac{\sinh \beta_m \sum_{k=1,3,\dots}^{\infty} \frac{1}{m^2} \left(\frac{m^2 + k^2}{a^2 + b^2} \right)^{-2} \left(\frac{2m^2 + k^2}{a^2 + b^2} \right) \left[\frac{k^2}{b^2} + (2-\mu) \frac{m^2}{a^2} \right]}{\frac{3+\mu}{1-\mu} \frac{\sinh(2\beta_m)}{2} + \beta_m} \times \right. \\ \left. \left[\left(\beta_m \coth \beta_m + 2 \frac{1+\mu}{1-\mu} \right) \times \sinh \frac{m\pi y}{a} - \frac{m\pi y}{a} \cosh \frac{m\pi y}{a} \right] \sin \frac{m\pi x}{a} \right\} \quad (33)$$

$$M_{y4} = -\frac{8M^T a}{\pi^2 b} \sum_{m=1,3,\dots}^{\infty} \left\{ \frac{\sinh \beta_m \sum_{k=1,3,\dots}^{\infty} \frac{1}{m^2} \left(\frac{m^2 + k^2}{a^2 + b^2} \right)^{-2} \left(\frac{2m^2 + k^2}{a^2 + b^2} \right) \left[\frac{k^2}{b^2} + (2-\mu) \frac{m^2}{a^2} \right]}{\frac{3+\mu}{1-\mu} \frac{\sinh(2\beta_m)}{2} + \beta_m} \times \right. \\ \left. \left(\frac{m\pi y}{a} \cosh \frac{m\pi y}{a} - \beta_m \coth \beta_m \sinh \frac{m\pi y}{a} \right) \sin \frac{m\pi x}{a} \right\} \quad (34)$$

叠加式(17)、式(28)、式(30)和式(32)得：

$$w = -\frac{4a^2 M^T}{D\pi^3} \sum_{m=1,3,\dots}^{\infty} \frac{1}{m^3 \cosh \alpha_m} \cosh \left(\frac{2\alpha_m y}{b} - \alpha_m \right) \sin \frac{m\pi x}{a} - \frac{M^T}{2D} (x-a)x + \\ \frac{2a^2(3-2\mu)M^T}{D\pi^3(1-\mu)} \sum_{m=1,3,\dots}^{\infty} \left\{ \frac{\sinh \beta_m \sinh \left(\frac{2\alpha_m y}{b} - \alpha_m \right)}{m^3 \cosh \alpha_m} \times \right. \\ \left. \frac{\frac{3+\mu}{1-\mu} \frac{\sinh(2\beta_m)}{2} + \beta_m}{\left[\left(\frac{2}{1-\mu} + \beta_m \coth \beta_m \right) \sinh \frac{m\pi y}{a} - \frac{m\pi y}{a} \cosh \frac{m\pi y}{a} \right] \sin \frac{m\pi x}{a}} \right\} - \\ \frac{8M^T}{\pi^4 D} \sum_{i=1,3,\dots}^{\infty} \sum_{j=1,3,\dots}^{\infty} \frac{1}{ij} \left(\frac{i^2}{a^2} + \frac{j^2}{b^2} \right)^{-2} \left(\frac{2i^2}{a^2} + \frac{j^2}{b^2} \right) \sin \frac{i\pi x}{a} \sin \frac{j\pi y}{b} - \\ \frac{8M^T a^3}{\pi^4 D b (1-\mu)} \sum_{m=1,3,\dots}^{\infty} \left\{ \frac{\sinh \beta_m \sum_{k=1,3,\dots}^{\infty} \frac{1}{m^4} \left(\frac{m^2 + k^2}{a^2 + b^2} \right)^{-2} \left(\frac{2m^2 + k^2}{a^2 + b^2} \right) \left[\frac{k^2}{b^2} + (2-\mu) \frac{m^2}{a^2} \right]}{\frac{3+\mu}{1-\mu} \frac{\sinh(2\beta_m)}{2} + \beta_m} \times \right. \\ \left. \left[\left(\frac{2}{1-\mu} + \beta_m \coth \beta_m \right) \sinh \frac{m\pi y}{a} - \frac{m\pi y}{a} \cosh \frac{m\pi y}{a} \right] \sin \frac{m\pi x}{a} \right\} \quad (35)$$

叠加式(18)、式(29)、式(31)和式(33), 得:

$$\begin{aligned}
 M_x = & \frac{4M^T}{\pi}(\mu-1) \sum_{m=1,3,\dots}^{\infty} \frac{1}{m \cosh \alpha_m} \cosh \left(\frac{m\pi y}{a} - \alpha_m \right) \sin \frac{m\pi x}{a} + \frac{2(3-2\mu)M^T}{\pi} \\
 & \sum_{m=1,3,\dots}^{\infty} \frac{\sinh \beta_m \sinh \left(\frac{2\alpha_m y}{b} - \alpha_m \right)}{m \cosh \alpha_m \left[\frac{3+\mu}{1-\mu} \frac{\sinh(2\beta_m)}{2} + \beta_m \right]} \left[\left(\beta_m \coth \beta_m + 2 \frac{1+\mu}{1-\mu} \right) \times \sinh \frac{m\pi y}{a} - \frac{m\pi y}{a} \cosh \frac{m\pi y}{a} \right] \sin \frac{m\pi x}{a} - \\
 & \frac{8M^T}{\pi^2} \sum_{i=1,3,\dots}^{\infty} \sum_{j=1,3,\dots}^{\infty} \frac{1}{ij} \left(\frac{i^2}{a^2} + \mu \frac{j^2}{b^2} \right) \left(\frac{i^2}{a^2} + \frac{j^2}{b^2} \right)^{-2} \left(\frac{2i^2}{a^2} + \frac{j^2}{b^2} \right) \sin \frac{i\pi x}{a} \sin \frac{j\pi y}{b} - \\
 & \frac{8M^T a}{\pi^2 b} \sum_{m=1,3,\dots}^{\infty} \left\{ \frac{\sinh \beta_m \sum_{k=1,3,\dots}^{\infty} \frac{1}{m^2} \left(\frac{m^2}{a^2} + \frac{k^2}{b^2} \right)^{-2} \left(\frac{2m^2}{a^2} + \frac{k^2}{b^2} \right) \left[\frac{k^2}{b^2} + (2-\mu) \frac{m^2}{a^2} \right]}{\frac{3+\mu}{1-\mu} \frac{\sinh(2\beta_m)}{2} + \beta_m} \times \right. \\
 & \left. \left[\left(\beta_m \coth \beta_m + 2 \frac{1+\mu}{1-\mu} \right) \times \sinh \frac{m\pi y}{a} - \frac{m\pi y}{a} \cosh \frac{m\pi y}{a} \right] \sin \frac{m\pi x}{a} \right\} \quad (36)
 \end{aligned}$$

叠加式(18)、式(29)、式(31)和式(34), 得:

$$\begin{aligned}
 M_y = & \frac{4M^T}{\pi}(1-\mu) \sum_{m=1,3,\dots}^{\infty} \frac{1}{m \cosh \alpha_m} \cosh \left(\frac{m\pi y}{a} - \alpha_m \right) \sin \frac{m\pi x}{a} + (\mu-1)M^T + \\
 & \frac{2(3-2\mu)M^T}{\pi} \sum_{m=1,3,\dots}^{\infty} \frac{\sinh \beta_m \sinh \left(\frac{2\alpha_m y}{b} - \alpha_m \right)}{m \cosh \alpha_m \left[\frac{3+\mu}{1-\mu} \frac{\sinh(2\beta_m)}{2} + \beta_m \right]} \left(\frac{m\pi y}{a} \cosh \frac{m\pi y}{a} - \beta_m \coth \beta_m \sinh \frac{m\pi y}{a} \right) \sin \frac{m\pi x}{a} - \\
 & \frac{8M^T}{\pi^2} \sum_{i=1,3,\dots}^{\infty} \sum_{j=1,3,\dots}^{\infty} \frac{1}{ij} \left(\mu \frac{i^2}{a^2} + \frac{j^2}{b^2} \right) \left(\frac{i^2}{a^2} + \frac{j^2}{b^2} \right)^{-2} \left(\frac{2i^2}{a^2} + \frac{j^2}{b^2} \right) \sin \frac{i\pi x}{a} \sin \frac{j\pi y}{b} - \\
 & \frac{8M^T a}{\pi^2 b} \sum_{m=1,3,\dots}^{\infty} \left\{ \frac{\sinh \beta_m \sum_{k=1,3,\dots}^{\infty} \frac{1}{m^2} \left(\frac{m^2}{a^2} + \frac{k^2}{b^2} \right)^{-2} \left(\frac{2m^2}{a^2} + \frac{k^2}{b^2} \right) \left[\frac{k^2}{b^2} + (2-\mu) \frac{m^2}{a^2} \right]}{\frac{3+\mu}{1-\mu} \frac{\sinh(2\beta_m)}{2} + \beta_m} \times \right. \\
 & \left. \left(\frac{m\pi y}{a} \cosh \frac{m\pi y}{a} - \beta_m \coth \beta_m \sinh \frac{m\pi y}{a} \right) \sin \frac{m\pi x}{a} \right\} \quad (37)
 \end{aligned}$$

3 结果分析

为了检验式(35)、式(36)和式(37)的正确性, 利用 MATLAB 软件编制程序进行计算。结果表明:

当取 $m=n=9$ 时, 挠度 w 已收敛于精确解; 对于单位宽度的弯矩 M_x , 当取 $m=n=17$ 时, 结果已收敛于精确解; 对于单位宽度的弯矩 M_y , 当取 $m=n=1999$

时,结果已基本收敛于精确解,此时与取 $m=n=2001$ 时的误差仅为 $1/10000$ 。为方便和工程实用起见,根据薄板的长宽比,将计算结果制成了表格(见表 2)。

表 2 三边固定一边自由矩形薄板壁面温差作用计算用表
Table 2 Bending calculation coefficient of three clamped sides and one free side under temperature disparity

$\mu = \frac{1}{6}$
 $M_x^T = k_x M^T$
 $M_y^T = k_y M^T$
 $M_{x1}^{OT} = k_{x1} M^T$
 $M_{y1}^{OT} = k_{y1} M^T$
 $M_{x2}^{OT} = k_{x2} M^T$
 $M_{y2}^{OT} = k_{y2} M^T$

$w(x, y) = f \frac{l_x^2 M^T}{D}$ (跨中挠度), 温度较低的一侧受拉。

l_x/l_y	k_{x1}	k_{y1}	k_{x2}	k_{y2}	k_x	k_y	f
0.50	1.0000	0.8333	0.8333	1.0000	0.2301	0.9699	0.0133
0.55	1.0000	0.8333	0.8333	1.0000	0.2350	0.9579	0.0191
0.60	1.0000	0.8333	0.8333	1.0000	0.2408	0.9398	0.0256
0.65	1.0000	0.8333	0.8333	1.0000	0.2437	0.9198	0.0326
0.70	1.0000	0.8333	0.8333	1.0000	0.2435	0.8985	0.0398
0.75	1.0000	0.8333	0.8333	1.0000	0.2403	0.8765	0.0470
0.80	1.0000	0.8333	0.8333	1.0000	0.2343	0.8542	0.0541
0.85	1.0000	0.8333	0.8333	1.0000	0.2260	0.8319	0.0608
0.90	1.0000	0.8333	0.8333	1.0000	0.2155	0.8099	0.0670
0.95	1.0000	0.8333	0.8333	1.0000	0.2031	0.7884	0.0726
1.00	1.0000	0.8333	0.8333	1.0000	0.1893	0.7676	0.0777
1.10	1.0000	0.8333	0.8333	1.0000	0.1584	0.7288	0.0858
1.20	1.0000	0.8333	0.8333	1.0000	0.1248	0.6935	0.0913
1.30	1.0000	0.8333	0.8333	1.0000	0.0905	0.6622	0.0943
1.40	1.0000	0.8333	0.8333	1.0000	0.0564	0.6345	0.0951
1.50	1.0000	0.8333	0.8333	1.0000	0.0234	0.6103	0.0941
1.60	1.0000	0.8333	0.8333	1.0000	-0.0077	0.5893	0.0915
1.70	1.0000	0.8333	0.8333	1.0000	-0.0368	0.5712	0.0878
1.80	1.0000	0.8333	0.8333	1.0000	-0.0638	0.5556	0.0832
1.90	1.0000	0.8333	0.8333	1.0000	-0.0885	0.5421	0.0779
2.00	1.0000	0.8333	0.8333	1.0000	-0.1113	0.5308	0.0722

为说明本文结论的工程实用性,以兰州地区某开敞式钢筋混凝土矩形贮液结构为例, C30 混凝土 ($E_c=3 \times 10^7 \text{kN/m}^2$), 尺寸为 $6\text{m} \times 6\text{m} \times 4.8\text{m}$, 壁板厚度为 250mm , 液体温度为 45°C 。考虑冬季最不利情况, 取室外计算温度为 -11°C , 则 $\alpha = 1 \times 10^{-5}/(^\circ\text{C})$, $\Delta T = 56^\circ\text{C}$ 。由表 2 $f=0.0541, k_x=0.2343, k_y=0.8542$ 。则 $w(3,2.4)=0.814\text{mm}$, $M_x(3,2.4)=24.6\text{kN} \cdot \text{m/m}$, $M_y(3,2.4)=89.7\text{kN} \cdot \text{m/m}$ 。由此可知, 尽管该结构在

液体温度 45°C 时所引起的跨中挠度较小(即 0.814mm), 但所引起的跨中弯矩很大(即 $M_x=24.6\text{kN} \cdot \text{m/m}, M_y=89.7\text{kN} \cdot \text{m/m}$)。设计中如不考虑由此所引起弯矩, 则在使用阶段会导致钢筋混凝土矩形贮液结构因配筋不足而开裂破坏。因此, 对于横向变温作用下的三边固支一边自由的混凝土矩形薄板(尤其对于跨度比较大的板结构), 为保证结构安全可靠、经济和正常使用, 应按本文提供的挠度及内力公式或表 2 进行强度和刚度的计算。

4 结论

本文基于薄板的小挠度理论和叠加原理, 推导出了横向变温作用下三边固支一边自由矩形薄板的解析解, 并利用 MATLAB 软件, 得到了基于混凝土材料的计算系数用表, 结果表明:

- (1) 弯矩和挠度计算公式的收敛性非常好, 从而证明了本文所得解析解的正确性。
- (2) 虽然横向变温作用的挠度较小, 可以忽略, 但是弯矩是不可忽略的, 设计计算时应引起足够地重视。
- (3) 所得计算系数用表方便实用, 工程设计时可按表 2 直接查表计算。

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